

How far do we understand our Universe at the moment?

Understanding Nature with the Spin-Charge-Family theory

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22st International Workshop "What comes beyond the standard
models?"

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Some publications:

- ▶ *Phys. Lett. B* **292**, 25-29 (1992), *J. Math. Phys.* **34**, 3731-3745 (1993), *Mod. Phys. Lett. A* **10**, 587-595 (1995), *Int. J. Theor. Phys.* **40**, 315-337 (2001),
- ▶ *Phys. Rev. D* **62** (04010-14) (2000), *Phys. Lett. B* **633** (2006) 771-775, **644** (2007) 198-202, **683** (2008) 110.1016, *JHEP* **04** (2014) 165, *Fortschritte Der Physik-Progress in Physics*, (2017) with H.B.Nielsen,
- ▶ *Phys. Rev. D* **74** 073013-16 (2006), with A.Borštnik Bračič,
- ▶ *New J. of Phys.* **10** (2008) 093002, arxiv:1412.5866, with G.Bregar, M.Breskvar, D.Lukman,
- ▶ *Phys. Rev. D* (2009) 80.083534, with G. Bregar,
- ▶ *New J. of Phys.* (2011) 103027, *J. Phys. A: Math. Theor.* **45** (2012) 465401, *J. Phys. A: Math. Theor.* **45** (2012) 465401, *J. of Mod. Phys.* **4** (2013) 823-847, arxiv:1409.4981, **6** (2015) 2244-2247, *Phys. Rev. D* **91** (2015) 6, 065004, . *J. Phys.: Conf. Ser. 845 01 IARD 2017*, *Eur. Phys. J.C.* **77** (2017) 231

More than **50 years ago** the **electroweak (and colour) standard model** offered an **elegant new step** in **understanding the origin of fermions and bosons** by **postulating**:

A.

- ▶ The existence of **massless family members** with the **charges** in the **fundamental representation of the groups** -
 - the **coloured triplet quarks and colourless leptons**,
 - the **left handed members as the weak charged doublets**,
 - the **right handed weak chargeless members**,
 - the **left handed quarks distinguishing in the hyper charge from the left handed leptons**,
 - **each right handed member having a different hyper charge**.
- ▶ The existence of **massless families to each of a family member**.

α name	hand- edness $-4iS^0 S^{12}$	weak charge τ^{13}	hyper charge Y	colour charge	elm charge Q
u_L^i	-1	$\frac{1}{2}$	$\frac{1}{6}$	colour triplet	$\frac{2}{3}$
d_L^i	-1	$-\frac{1}{2}$	$\frac{1}{6}$	colour triplet	$-\frac{1}{3}$
ν_L^i	-1	$\frac{1}{2}$	$-\frac{1}{2}$	colourless	0
e_L^i	-1	$-\frac{1}{2}$	$-\frac{1}{2}$	colourless	-1
u_R^i	1	weakless	$\frac{2}{3}$	colour triplet	$\frac{2}{3}$
d_R^i	1	weakless	$-\frac{1}{3}$	colour triplet	$-\frac{1}{3}$
ν_R^i	1	weakless	0	colourless	0
e_R^i	1	weakless	-1	colourless	-1

Members of each of the $i = 1, 2, 3$ families, $i = 1, 2, 3$ massless before the electroweak break. Each family contains the left handed weak charged quarks and the right handed weak chargeless quarks, belonging to the colour triplet $(1/2, 1/(2\sqrt{3}))$, $(-1/2, 1/(2\sqrt{3}))$, $(0, -1/(\sqrt{3}))$.

And the anti-fermions to each family and family member.

B.

- ▶ The existence of **massless vector gauge fields** to the observed **charges** of the **family members**, **carrying charges** in the **adjoint representation of the charge groups**.

Gauge fields before the electroweak break

- ▶ Three massless vector fields, the gauge fields of the three charges.

name	hand- edness	weak charge	hyper charge	colour charge	elm charge
hyper photon	0	0	0	colourless	0
weak bosons	0	triplet	0	colourless	triplet
gluons	0	0	0	colour octet	0

They all are vectors in $d = (3 + 1)$, in the adjoint representations with respect to the weak, colour and hyper charges.

C.

- ▶ The **existence of a massive scalar field - the higgs**,
 - carrying the weak charge $\pm\frac{1}{2}$ and the hyper charge $\mp\frac{1}{2}$ — as it would be in the **fundamental representation** of the **groups**,
 - gaining at some step a **"nonzero vacuum expectation values"**, breaking the weak and the hyper charge and correspondingly breaking the **mass protection**.
- ▶ The **existence** of the **Yukawa couplings**, taking care of
 - the properties of **fermions** and
 - the masses of the **heavy bosons**.

- ▶ The Higgs's field, the scalar in $d = (3 + 1)$, a doublet with respect to the weak charge.

name	hand- edness	weak charge	hyper charge	colour charge	elm charge
0. $Higgs_u$	0	$\frac{1}{2}$	$\frac{1}{2}$	colourless	1
$\langle Higgs_d \rangle$	0	$-\frac{1}{2}$	$\frac{1}{2}$	colourless	0

name	hand- edness	weak charge	hyper charge	colour charge	elm charge
$\langle Higgs_u \rangle$	0	$\frac{1}{2}$	$-\frac{1}{2}$	colourless	0
0. $Higgs_d$	0	$-\frac{1}{2}$	$-\frac{1}{2}$	colourless	-1

D.

- ▶ There is the **gravitational field** in **$d=(3+1)$** .

- ▶ **The *standard model* assumptions have been confirmed without offering surprises.**
- ▶ The last unobserved field as a field, the **Higgs's scalar**, detected in June 2012, was confirmed in March 2013.
- ▶ The waves of the **gravitational field** were detected in February 2016 and again 2017.

- ▶ The *standard model* assumptions have in the literature several explanations, but with **many new not explained assumptions**.
- ▶ I am proposing **the *spin-charge-family* theory**, which offers the explanation for
 - all the assumptions of the *standard model*,
 - for many observed phenomena,
 - making several predictions.**

Is this the right next step beyond both *standard models*?

There are namely many phenomena

- ▶ the **dark matter**,
 - ▶ the **matter-antimatter** asymmetry,
 - ▶ the **dark energy**,
 - ▶ the **observed dimension of space time**,
 - ▶ **many other phenomena, not yet understood.**
-
- ▶ Can the *spin-charge-family theory* explain the observed phenomena explaining as well all the assumptions of the *standard model*?

Work done so far on the *spin-charge-family theory* is very promising.

To better understand what is essential in physics beyond the standard models two projects are proposed:

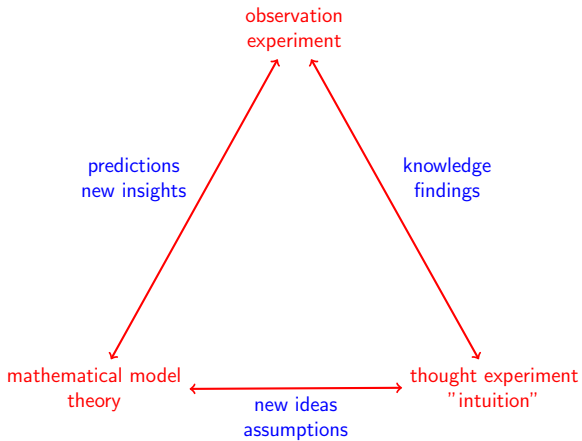
- ▶ **Could exist fermions with integer spins, that is with spins and charges in the fundamental representations of the group in $d \geq 5$ — like in the *spin-charge family theory*?**
N.S.M.B. and H.B.N. and..
- ▶ **Can one bosonize fermions or fermionize bosons?**
H.B.N. and N.S.M.B. and..

How do we gather knowledge - understanding -

- ▶ **about the laws, which govern the smallest constituents of Nature**
- ▶ **and correspondingly the whole our Universe and might be universes?**
- ▶ **How do we learn about the space-time in which we live?**

- ▶ We observe.
- ▶ We make experiments.
- ▶ We make thought experiments, leading to new questions, new ideas.
- ▶ We make mathematical models, which cover the so far made discoveries, findings, thoughts, ideas.
- ▶ The models make predictions.
- ▶ We check predictions by experiments.

We repeat this circle again and again, with more and more knowledge and more and more and more intuition.



We try to understand:

- ▶ What are elementary constituents and interactions among constituents in our Universe, in any universe?
- ▶ Can the elementary constituent be of only one kind? Are the four observed interactions — gravitational, elektromagnetic, weak and colour — of the common origin?
- ▶ Is the space-time the so far observed $(3 + 1)$? Why $(3+1)$?
- ▶ If not $(3 + 1)$ may it be that the space-time is infinite?
- ▶ How has the space-time of our universe started?
- ▶ What is the matter and what the anti-matter?

Obviously it is the time to make **the next step beyond both standard models.**

What questions should one ask to be able to find **next steps** beyond the *standard models* and to understand not yet understood phenomena?

- ▶ ○ Where do **family members** originate?
 - Where do **charges** of **family members** originate?
 - Why are the **charges** of **family members** so different?
 - Why have the **left handed family members** so different charges from the **right handed** ones?
- ▶ ○ Where do **families** of **family members** originate?
 - How **many different families** exist?
 - Why do **family members – quarks and leptons** – manifest so different properties if they all start as massless?

- ▶ **o** How is the **origin** of the **scalar field** (the Higgs's scalar) and the **Yukawa couplings connected** with the origin of **families**?
 - o** How many **scalar fields** determine properties of the so far (and others possibly be) **observed fermions** and masses of **weak bosons**? (The Yukawa couplings certainly speak for the existence of several scalar fields with the properties of Higgs's scalar.)
- ▶ **Why is the Higgs's scalar**, or are all **scalar fields**, if there are several, **doublets** with respect to the weak and the hyper charge?
- ▶ **Do exist** also **scalar fields** with the **colour charge in the fundamental representation** and where, if they are, **do they manifest**?

- ▶ Where do the **charges** and correspondingly the so far (and others possibly be) **observed vector gauge fields** originate?
- ▶ **Where** does the **dark matter** originate?
- ▶ **Where** does the "ordinary" **matter-antimatter asymmetry** originate?
- ▶ **Where** does the **dark energy** originate?
- ▶ What is the dimension of space? $(3 + 1)?$, $((d - 1) + 1)?$, $\infty?$
- ▶ **What** is the role of the **symmetries**– discrete, continuous, global and gauge – in our **universe, in Nature?**
- ▶ And many others.

My statement:

- ▶ **An elegant trustworthy next step** must offer answers to **several** open questions, explaining:
 - o The **origin of the family members and the charges.**
 - o The **origin of the families and their properties.**
 - o The **origin of the scalar fields and their properties.**
 - o The **origin of the vector fields and their properties.**
 - o The **origin of the dark matter.**
 - o The **origin of the "ordinary" matter-antimatter asymmetry.**

My statement continues:

- ▶ There exist not yet observed families, gauge vector and scalar gauge fields.
- ▶ **Dimension of space is larger than 4** (very probably infinite).
- ▶ Inventing a next step which covers one of the open questions, might be of a help **but can hardly show the right next step in understanding nature.**

In the literature **NO explanation for the existence of the families can be found**, which would not just assume the family groups.

Several extensions of the **standard model** are, however, proposed, like:

- ▶ The $SU(3)$ group is assumed to describe – not explain – the existence of three families.

Like the **Higgs's** scalar charges are in the **fundamental** representations of the groups, also the **Yukawas** are assumed to emerge from the scalar fields, in the **fundamental** representation of the $SU(3)$ group.

- ▶ **SU(5) and SO(10) grand unified theories are proposed, unifying all the charges.** But the **spin** (the handedness) is obviously connected with the (weak and the hyper) charges, what these theories do "by hand" as it does the *standard model*, and the appearance of families is not explained.
- ▶ **Supersymmetric theories, assuming the existence of bosons with the charges of quarks and leptons and fermions with the charges of the gauge vector fields, although having several nice properties but not explaining the appearance of families (except again by assuming larger groups), are not, to my understanding, the right next step beyond the *standard model*.**

- o The **Spin-Charge-Family** theory does offer the **explanation for all the assumptions of the standard model**, answering up to now several of the above cited open questions!
 - o The **more effort** is put into this theory, the more answers to the open questions in elementary particle physics and cosmology is the theory offering.

- o I shall first make a short introduction into the **Spin-Charge-Family** theory.
- o I shall make an overview of achievements so far of the **Spin-Charge-Family** theory.
- o I shall report on the possibility that "**nature could make a choice**" of **Grassmann** rather than **Clifford space** to describe the internal degrees of freedom of fermions, what would lead to fermions, still anticommuting (in the second quantization), but with the integer spin [arXiv:1802.05554v1v2].

- ▶ A brief introduction into the **spin-charge-family theory**.

- ▶ **Spinors** carry in $d \geq (13 + 1)$ **two kinds of spin, no charges**, *Phys. Rev. D* **91** 065004 (2015), *J.of Mod. Physics* **6** (2015) 2244.
 - The **Dirac spin** (γ^a) in $d = (13 + 1)$ **describes** in $d = (3 + 1)$ **spin** and **ALL the charges of quarks and leptons, left and right handed**, explaining all the assumptions about the charges and the handedness of the standard model *J. of Math. Phys.* **34** (1993), 3731, *J. of Math. Phys.* **43**, 5782 (2002) [hep-th/0111257].
 - The **second kind of spin** ($\tilde{\gamma}^a$) **describes FAMILIES**, explaining the origin of families and their number, *J. of Math. Phys.* **44** 4817 (2003) [hep-th/0303224].
 - There is **NO third kind of spin**.
- ▶ **C,P,T symmetries** in $d = (3 + 1)$ **follow from the C,P,T symmetry** in $d \geq (13 + 1)$. (*JHEP* **04** (2014) 165)

- ▶ All **vector** and **scalar gauge fields** origin in gravity, explaining the origin of the vector and scalar gauge fields, which in the *standard model* are assumed *Eur. Phys. J. C* **77** (2017) 231:

- o in two **spin connection fields**, the **gauge fields** of $\gamma^a \gamma^b$ and $\tilde{\gamma}^a \tilde{\gamma}^b$, and in

- o **vielbeins**, the gauge fields of moments *Eur. Phys. J. C* **77** (2017) 231, [arXiv:1604.00675]

- ▶ If there are **no spinor sources present**, then either vector (\vec{A}_m^A , $m = 0, 1, 2, 3$) or scalar (\vec{A}_s^A , $s = 5, 6, \dots, d$) gauge fields are determined by **vielbeins** uniquely.

- ▶ **Spinors** interact correspondingly with
 - the **vielbeins** and
 - the **two kinds of the spin connection fields**, *Eur. Phys. J. C* **77** (2017) 231.
- ▶ In $d = (3 + 1)$ the **spin-connection fields**, together with the **vielbeins**, manifest either as
 - **vector** gauge fields with all the **charges** in the **adjoint** representations or as
 - **scalar** gauge fields with the **charges** with respect to the **space index** in the **"fundamental"** representations and all the other **charges** in the **adjoint** representations or as
 - **tensor** gravitational field.

There are two kinds of **scalar fields** with respect to the space index s :

- ▶ Those with ($s = 5, 6, 7, 8$) (they carry zero "**spinor charge**") are **doublets** with respect to the $SU(2)_I$ (the weak) charge and the **second $SU(2)_{II}$ charge** (determining the hyper charge). They are in the **adjoint** representations with respect to the **family** and the **family members charges**.
- These **scalars** explain the **Higgs's scalar** and the **Yukawa couplings**.

Phys. Rev. **D 91** (2015) 6, 065004

- ▶ **Those** with twice the "spinor charge" of a quark and ($s = 9, 10, ..d$) are **colour triplets**. Also they are in the **adjoint representations** with respect to the **family** and the **family members charges**.
 - These **scalars** transform **antileptons** into **quarks**, and **antiquarks** into **quarks** and back and correspondingly **contribute to matter-antimatter asymmetry** of our universe and to **proton decay**.
- ▶ There are **no additional scalar fields** in the **spin-charge-family theory**, if $d = (13 + 1)$.

Phys. Rev. **D 91** (2015) 6, 065004

J. of Mod. Phys. **6** (2015) 2244

Condensate

- ▶ The (assumed so far, waiting to be derived how does this spontaneously appear) **scalar condensate** of **two right handed neutrinos** with the **family** quantum numbers of the upper four families (there are two four family groups in the theory), appearing $\approx 10^{16}$ GeV or higher,
 - **breaks the CP** symmetry, causing the **matter-antimatter asymmetry** and the proton decay,
 - couples to all the **scalar fields**, making them massive,
 - couples to all the phenomenologically **unobserved vector gauge fields**, making them massive.
 - Before the electroweak break all the so far **observed vector gauge fields are massless**.

Phys. Rev. **D 91** (2015) 6, 065004,

J. of Mod. Phys. **6** (2015) 2244,

J. Phys.: Conf.Ser. 845 01, **IARD 2017**

- ▶ The **vector fields**, which do not couple to the condensate and remain massless, are:
 - o the **hyper charge vector field**.
 - o the **weak vector fields**,
 - o the **colour vector fields**,
 - o the **gravity fields**.

The $SU(2)_{II}$ symmetry breaks due to the **condensate**, leaving the **hyper charge unbroken**.

Nonzero vacuum expectation values of scalars

- waiting to be shown how does such an event appear in the *spin-charge-family* spontaneously.
- ▶ The scalar fields with the **space index (7, 8)**, gaining **nonzero vacuum expectation values**, cause the **electroweak break**,
 - breaking the weak and the hyper charge,
 - changing their own masses,
 - bringing masses to the **weak bosons**,
 - bringing masses to the **families of quarks and leptons**.

Phys. Rev. **D 91** (2015) 6, 065004,
J. Phys.: Conf.Ser. 845 01 **IARD 2017**,
Eur. Phys. J.C. **77** (2017) 231 [arXiv:1604.00675],
J. of Mod. Phys. **6** (2015) 2244, [arXiv:1502.06786,
arXiv:1409.4981]

- ▶ The only gauge fields which do not couple to these scalars and remain massless are
 - electromagnetic,
 - colour vector gauge fields,
 - gravity.
- ▶ There are two times four decoupled massive **families** of **quarks and leptons** after the electroweak break:
 - There are the observed **three families** among the **lower four, the fourth to be observed**.
 - The stable among the **upper four families** form the **dark matter**.

Phys. Rev. **D 80**, 083534 (2009),

Phys. Rev. **D 91** (2015) 6, 065004,

J. Phys.: Conf.Ser. 845 01, **IARD 2017**

- ▶ All the families are **singlets** with respect to $\widetilde{SU}(3)$ group, originating in the second kind of the Clifford algebra object $\tilde{\gamma}^a \tilde{\gamma}^b$.

Phys. Lett. **B** 633 (2006) 771-775 [hep-th/0509101]
arXiv:0812.0510[hep-ph],44-55

- ▶ It is **extremely encouraging** for the **spin-charge-family theory**, that a **simple starting action** contains **all the degrees of freedom observed at low energies**, directly or indirectly, and that only the
 - **condensate** and
 - **nonzero vacuum expectation values of all the scalar fields with $s = (7, 8)$**are needed that the **theory explains**
 - **all the assumptions** of the standard model, with the gauge fields, scalar fields, families of fermions, masses of fermions and of bosons included,
 - explaining also **the dark matter**,
 - **the matter/antimatter asymmetry**,
 - **the triangle anomalies cancellation** in the standard model (Forts. der Physik, Prog.of Phys.) (2017) 1700046) and...

The **spin-charge-family** theory is a kind of a **Kaluza-Klein-like** theory, but with **two kinds of spins**.

In d -dimensional space there are **fermions** with **two kinds of spins** and **gravity**, represented by two **spin connection** and **vielbein gauge fields**.

J. of Mod. Phys. **4** (2013) 823,
Phys. Rev. **D 91** (2015) 6, 065004,
J. of Mod. Phys. **6** (2015) 2244 [arxiv:1409.4981],
IARD, J. Phys.: Conf. Ser. 845 (2017) 012017
[arXiv:1607.01618]/v2],
Eur. Phys. J.C. **77** (2017) 231,
Forts. der Physik, Prog. of Phys.) (2017) 1700046

Comparing the **spin-charge-family** and the **unifying theories**, with the **Kaluza-Klein like** theories included.

- ▶ The $SO(10)$ must be put into $SO(13, 1)$,
- ▶ allowing **only gravity** as **gauge fields**,
- ▶ recognizing that there are **two kinds of the Clifford algebra** objects
 - 0 one **explaining spins and charges** of fermions,
 - 0 the second one **explaining families** of fermions.

This would be the first step towards the **spin-charge-family**, if assuming as well the **simple starting action for fermions** and **gauge fields**, which at the low energies offers all the observed degrees of freedom, so that nothing in addition is needed.

- ▶ One has to recognize:
- ▶ that the **scalar gauge fields** are of **two kinds** due to the two **kinds of the Clifford algebra** objects,
- ▶ offering the explanation for several so far observed phenomena,
 - 0 for the masses of fermions and weak bosons,
 - 0 for the appearance of the matter-antimatter asymmetry in the universe,
 - 0 for the appearance of the dark matter,
 - 0 and for several other phenomena.
- ▶ One must also demonstrate that breaking symmetry can still lead to (almost) massless fermions,
- ▶ proving this, making several predictions.

- ▶ In addition, as I shall discuss in this talk, the **theory offers the second quantized fields without assumptions**, what the ordinary theories, assuming the needed groups to represent spins, charges and families, must make.

A short look "inside" the **spin-charge-family** theory.

There are **only two kinds of the Clifford algebra objects in any d**:

- ▶ The **Dirac γ^a operators** (used by Dirac 90 years ago).
- ▶ The **second one: $\tilde{\gamma}^a$** , which I recognized in the Grassmann space.

$$\{\gamma^a, \gamma^b\}_+ = 2\eta^{ab} = \{\tilde{\gamma}^a, \tilde{\gamma}^b\}_+,$$

$$\{\gamma^a, \tilde{\gamma}^b\}_+ = 0,$$

$$(\tilde{\gamma}^a \mathbf{B} : = \mathbf{i}(-)^{n_B} \mathbf{B} \gamma^a) |\psi_0 \rangle,$$

$$(\mathbf{B} = a_0 + a_a \gamma^a + a_{ab} \gamma^a \gamma^b + \cdots + a_{a_1 \dots a_d} \gamma^{a_1} \dots \gamma^{a_d}) |\psi_0 \rangle$$

$(-)^{n_B} = +1, -1$, when the object B has a Clifford even or odd character, respectively.

$|\psi_0 \rangle$ is a vacuum state on which the operators γ^a **apply**.

$$\mathbf{S}^{ab} := (\mathbf{i}/4)(\gamma^a \gamma^b - \gamma^b \gamma^a),$$

$$\tilde{\mathbf{S}}^{ab} := (\mathbf{i}/4)(\tilde{\gamma}^a \tilde{\gamma}^b - \tilde{\gamma}^b \tilde{\gamma}^a),$$

$$\{\mathbf{S}^{ab}, \tilde{\mathbf{S}}^{cd}\}_- = \mathbf{0}.$$

- ▶ $\tilde{\mathbf{S}}^{ab}$ define the equivalent representations with respect to \mathbf{S}^{ab} .

My recognition:

- ▶ If γ^a are used to describe **the spin and the charges of spinors**,
 $\tilde{\gamma}^a$ - since it must be used or it must be explained why it does not manifest - it **must be used to describe families of spinors**.

Phys. Lett. **B 292**, 25-29 (1992),

J. Math. Phys. **34**, 3731-3745 (1993),

Mod. Phys. Lett. **A 10**, 587-595 (1995)

A simple action for a **spinor** which carries in $d = (13 + 1)$ only **two kinds of spins** (no charges) and for **gauge fields**:

$$\mathbf{S} = \int d^d x E \mathcal{L}_f + \int d^d x E (\alpha R + \tilde{\alpha} \tilde{R})$$



$$\begin{aligned} \mathcal{L}_f &= \frac{1}{2}(\bar{\psi} \gamma^a p_{0a} \psi) + h.c. \\ p_{0a} &= f^\alpha{}_a p_{0\alpha} + \frac{1}{2E} \{p_\alpha, E f^\alpha{}_a\} - \\ \mathbf{p}_{0\alpha} &= \mathbf{p}_\alpha - \frac{1}{2} \mathbf{S}^{ab} \omega_{ab\alpha} - \frac{1}{2} \tilde{\mathbf{S}}^{ab} \tilde{\omega}_{ab\alpha} \end{aligned}$$

- ▶ The Einstein action for a free gravitational field is assumed to be linear in the curvature

$$\begin{aligned}
 \mathcal{L}_g &= E (\alpha \mathbf{R} + \tilde{\alpha} \tilde{\mathbf{R}}), \\
 \mathbf{R} &= f^\alpha [a f^{\beta b}] (\omega_{ab\alpha,\beta} - \omega_{ca\alpha} \omega^c_{b\beta}), \\
 \tilde{\mathbf{R}} &= f^\alpha [a f^{\beta b}] (\tilde{\omega}_{ab\alpha,\beta} - \tilde{\omega}_{ca\alpha} \tilde{\omega}^c_{b\beta}),
 \end{aligned}$$

with $E = \det(e^a_\alpha)$
 and $f^\alpha [a f^{\beta b}] = f^{\alpha a} f^{\beta b} - f^{\alpha b} f^{\beta a}$.

- ▶ The only internal degrees of freedom of **spinors** (**fermions**) are the **two kinds of the spin**.
- ▶ The only **gauge fields** are the **gravitational** ones – **vielbeins and the two kinds of spin connections**.
- ▶ Either γ^a or $\tilde{\gamma}^a$ transform as vectors in d ,
and correspondingly also $f^\alpha{}_a \omega_{bc\alpha}$ and $f^\alpha{}_a \tilde{\omega}_{bc\alpha}$ transform as tensors with respect to the flat index a .

Variation of the action brings for $\omega_{ab\alpha}$

$$\begin{aligned}
 \omega_{ab\alpha} = & -\frac{1}{2E} \left\{ e_{e\alpha} e_{b\gamma} \partial_\beta (E f^{\gamma[e} f^{\beta]}_{a]}) + e_{e\alpha} e_{a\gamma} \partial_\beta (E f^{\gamma}{}_{[b} f^{\beta e]}) \right. \\
 & \left. - e_{e\alpha} e^e{}_\gamma \partial_\beta (E f^{\gamma}{}_{[a} f^{\beta]}_{b]}) \right\} \\
 & - \frac{e_{e\alpha}}{4} \left\{ \bar{\Psi} \left(\gamma_e S_{ab} + \frac{3i}{2} (\delta_b^e \gamma_a - \delta_a^e \gamma_b) \right) \Psi \right\} \\
 & - \frac{1}{d-2} \left\{ e_{a\alpha} \left[\frac{1}{E} e^d{}_\gamma \partial_\beta (E f^{\gamma}{}_{[d} f^{\beta]}_{b]}) + \frac{1}{2} \bar{\Psi} \gamma^d S_{db} \Psi \right] \right. \\
 & \left. - e_{b\alpha} \left[\frac{1}{E} e^d{}_\gamma \partial_\beta (E f^{\gamma}{}_{[d} f^{\beta]}_{a]}) + \frac{1}{2} \bar{\Psi} \gamma^d S_{da} \Psi \right] \right\}
 \end{aligned}$$

IARD, J. Phys.: Conf. Ser. 845 012017 and the refs. therein

and for $\tilde{\omega}_{ab\alpha}$,

$$\begin{aligned}
 \tilde{\omega}_{ab\alpha} = & -\frac{1}{2E} \left\{ e_{e\alpha} e_{b\gamma} \partial_\beta (E f^{\gamma[e} f^{\beta]}_{a]}) + e_{e\alpha} e_{a\gamma} \partial_\beta (E f^{\gamma}_{[b} f^{\beta e]}) \right. \\
 & \left. - e_{e\alpha} e^e{}_\gamma \partial_\beta (E f^{\gamma}_{[a} f^{\beta]}_{b]}) \right\} \\
 & - \frac{e_{e\alpha}}{4} \left\{ \bar{\Psi} \left(\gamma_e \tilde{S}_{ab} + \frac{3i}{2} (\delta_b^e \gamma_a - \delta_a^e \gamma_b) \right) \Psi \right\} \\
 & - \frac{1}{d-2} \left\{ e_{a\alpha} \left[\frac{1}{E} e^d{}_\gamma \partial_\beta (E f^{\gamma}_{[d} f^{\beta]}_{b]}) + \frac{1}{2} \bar{\Psi} \gamma^d \tilde{S}_{db} \Psi \right] \right. \\
 & \left. - e_{b\alpha} \left[\frac{1}{E} e^d{}_\gamma \partial_\beta (E f^{\gamma}_{[d} f^{\beta]}_{a]}) + \frac{1}{2} \bar{\Psi} \gamma^d \tilde{S}_{da} \Psi \right] \right\}
 \end{aligned}$$

Eur. Phys. J. C, **77** (2017) 231 and the refs. therein

Fermions

- ▶ The action for **spinors seen** from $d = (3 + 1)$ and **analyzed** with respect to the standard model groups as subgroups of $SO(1 + 13)$:

$$\begin{aligned}
 \mathcal{L}_f = & \bar{\psi} \gamma^m (p_m - \sum_{A,i} g^{A\tau} A_i^{Ai}) \psi + \\
 & \left\{ \sum_{s=[7],[8]} \bar{\psi} \gamma^s p_{0s} \psi \right\} + \\
 & \left\{ \sum_{s=[5],[6]} \bar{\psi} \gamma^s p_{0s} \psi + \right. \\
 & \left. \sum_{t=[9],\dots,[14]} \bar{\psi} \gamma^t p_{0t} \psi \right\} \cdot \\
 & + \text{the rest } , ,
 \end{aligned}$$

Covariant momenta

$$p_{0m} = \left\{ p_m - \sum_A g^A \vec{\tau}^A \vec{A}_m^A \right\}$$

$$\mathbf{m} \quad n \quad (0, 1, 2, 3),$$

$$p_{0s} = f_s^\sigma \left[p_\sigma - \sum_A g^A \vec{\tau}^A \vec{A}_\sigma^A - \sum_A \tilde{g}^A \vec{\tau}^A \vec{A}_\sigma^A \right],$$

$$\mathbf{s} \in (7, 8),$$

$$p_{0s} = f_s^\sigma \left[p_\sigma - \sum_A g^A \vec{\tau}^A \vec{A}_\sigma^A - \sum_A \tilde{g}^A \vec{\tau}^A \vec{A}_\sigma^A \right],$$

$$\mathbf{s} \in (5, 6),$$

$$p_{0t} = f_t^{\sigma'} \left(p_{\sigma'} - \sum_A g^A \vec{\tau}^A \vec{A}_{\sigma'}^A - \sum_A \tilde{g}^A \vec{\tau}^A \vec{A}_{\sigma'}^A \right),$$

$$\mathbf{t} \in (9, 10, 11, \dots, 14),$$

$$\mathbf{A}_s^{\text{Ai}} = \sum_{a,b} \mathbf{c}_{ab}^{\text{Ai}} \omega_{\text{abs}} ,$$

$$\mathbf{A}_t^{\text{Ai}} = \sum_{a,b} \mathbf{c}_{ab}^{\text{Ai}} \omega_{\text{abt}} ,$$

$$\tilde{\mathbf{A}}_s^{\text{Ai}} = \sum_{a,b} \tilde{\mathbf{c}}_{ab}^{\text{Ai}} \tilde{\omega}_{\text{abs}} ,$$

$$\tilde{\mathbf{A}}_t^{\text{Ai}} = \sum_{a,b} \tilde{\mathbf{c}}_{ab}^{\text{Ai}} \tilde{\omega}_{\text{abt}} .$$

$$\tau^{Ai} = \sum_{a,b} c^{Ai}_{ab} \mathbf{S}^{ab},$$

$$\tilde{\tau}^{Ai} = \sum_{a,b} \tilde{c}^{Ai}_{ab} \tilde{\mathbf{S}}^{ab},$$

$$\{\tau^{Ai}, \tau^{Bj}\}_- = i\delta^{AB} f^{Aijk} \tau^{Ak},$$

$$\{\tilde{\tau}^{Ai}, \tilde{\tau}^{Bj}\}_- = i\delta^{AB} f^{Aijk} \tilde{\tau}^{Ak},$$

$$\{\tau^{Ai}, \tilde{\tau}^{Bj}\}_- = 0.$$

- ▶ τ^{Ai} represent the *standard model* charge groups
— $SU(3)_c, SU(2)_w$ — the second $SU(2)_{II}$, the "spinor"
charge $U(1)$, taking care of the hyper charge Y ,
- ▶ $\tilde{\tau}^{Ai}$ denote the family quantum numbers.

$$\mathbf{N}_{(L,R)}^i := \frac{1}{2}(S^{23} \pm iS^{01}, S^{31} \pm iS^{02}, S^{12} \pm iS^{03}),$$

$$\tau_{(1,2)}^i := \frac{1}{2}(S^{58} \mp S^{67}, S^{57} \pm S^{68}, S^{56} \mp S^{78}),$$

$$\tau_3^i := \frac{1}{2}\{S^{9\ 12} - S^{10\ 11}, S^{9\ 11} + S^{10\ 12}, S^{9\ 10} - S^{11\ 12}, \\ S^{9\ 14} - S^{10\ 13}, S^{9\ 13} + S^{10\ 14}, S^{11\ 14} - S^{12\ 13}, \\ S^{11\ 13} + S^{12\ 14}, \frac{1}{\sqrt{3}}(S^{9\ 10} + S^{11\ 12} - 2S^{13\ 14})\},$$

$$\tau^4 := -\frac{1}{3}(S^{9\ 10} + S^{11\ 12} + S^{13\ 14}),$$

$$Y := \tau^4 + \tau^{23},$$

$$Y' := -\tau^4 \tan^2 \vartheta_2 + \tau^{23},$$

$$Q := \tau^{13} + Y,$$

$$Q' := -Y \tan^2 \vartheta_1 + \tau^{13},$$

and equivalently for family groups \tilde{S}^{ab} .

Breaks of symmetries after starting with

- o massless spinors (fermions) ,
- o vielbeins and two kinds of the spin connection fields

**We prove for a toy model that breaking symmetry in
Kaluza-Klein theories can lead to massless fermions.**

New J. Phys. 13, 103027, 2011.

J. Phys. A: Math. Theor. 45, 465401, 2012.

[arXiv:1205.1714], [arXiv:1312.541], [arXiv:hep-ph/0412208
p.64-84].

[arXiv:1302.4305], p. 157-166.

$$SO(1, 13) \times \widetilde{SO(1, 13)} \times$$

BREAK I
 at $E \geq 10^{16} \text{ GeV}$



$$SO(1, 7) \times$$

$$\times \widetilde{SO(1, 7)}$$

$$U(1) \times$$

$$SU(3)$$



eight massless families



$$SO(1, 3) \times SO(4) \times U(1) \times$$

$$(\widetilde{SU(2)}_{I_{\widetilde{SO(1,3)}}} \times \widetilde{SU(2)}_{II_{\widetilde{SO(4)}}}) \times$$

(divided into two groups)

$$(\widetilde{SU(2)}_{II_{\widetilde{SO(1,3)}}} \times \widetilde{SU(2)}_{II_{\widetilde{SO(4)}}}) \times SU(3)$$

BREAK II



The Standard Model like way of breaking



$$SO(1, 3) \times U(1) \times SU(3)$$

× (two groups of four massive families)

- ▶ Both breaks leave **eight families** ($2^{8/2-1} = 8$, determined by the symmetry of $\widetilde{SO}(1,7)$) massless. All the families are $\widetilde{SU}(3)$ **chargeless**.

Phys. Rev. **D**, 80.083534 (2009)

- ▶ The appearance of the **condensate of the two right handed neutrinos**, coupled to **spin 0**, makes the boson gauge fields, with which the condensate interacts, massive. These gauge fields are:

- All the scalar gauge fields with the space index $s \geq 5$.

- The vector ($m \leq 3$) gauge fields with the Y' charges — the superposition of $SU(2)_{//}$ and $U(1)_{//}$ charges.

J. Phys.: Conf. Ser. 845 (2017) 012017

The **condensate** has spin $S^{12} = 0$, $S^{03} = 0$,
 weak charge $\vec{\tau}^1 = 0$, and
 $\vec{\tilde{\tau}}^1 = 0$, $\tilde{Y} = 0$, $\tilde{Q} = 0$, $\vec{\tilde{N}}_L = 0$.

state	τ^{23}	τ^4	Y	Q	$\tilde{\tau}^{23}$	\tilde{N}_R^3	$\tilde{\tau}^4$
$ \nu_{1R}^{VIII} \rangle_1 \nu_{2R}^{VIII} \rangle_2$	1	-1	0	0	1	1	-1
$ \nu_{1R}^{VIII} \rangle_1 e_{2R}^{VIII} \rangle_2$	0	-1	-1	-1	1	1	-1
$ e_{1R}^{VIII} \rangle_1 e_{2R}^{VIII} \rangle_2$	-1	-1	-2	-2	1	1	-1

- ▶ The **colour, elm, weak and hyper** vector gauge fields do **not interact with the condensate and remain massless.**
J. of Mod. Physics **6** (2015) 2244

- ▶ **At the electroweak break** from $SO(1,3) \times SU(2)_I \times U(1)_I \times SU(3)$ to $SO(1,3) \times U(1) \times SU(3)$
 - scalar fields with the space index $s = (7, 8)$ obtain nonzero vacuum expectation values,
 - break correspondingly the weak and the hyper charge and change their own masses.
 - They leave massless only the **colour, elm** and **gravity gauge fields**.
- ▶ All the eight massless families gain masses.

Also these is so far just assumed, waiting to be proven that scalar fields, together with boundary conditions, are spontaneously causing also this last breaks.

But all the needed degrees of freedom are already in the simple starting action, looking so far very promising.

- ▶ To the **electroweak break** several scalar fields, the gauge fields of **twice $\widetilde{SU}(2) \times \widetilde{SU}(2)$ and three $\times U(1)$** , contribute, all with the **weak and the hyper charge** of the *standard model Higgs*.
- ▶ They carry besides the **weak** and the **hyper charge** either
 - o the **family members** quantum numbers originating in **(Q,Q',Y')** or
 - o the **family** quantum numbers originating in **twice $\widetilde{SU}(2) \times \widetilde{SU}(2)$** .

J. of Mod. Physics **6** (2015) 2244.

- ▶ **The mass matrices of each family member manifest the $\widetilde{SU}(2) \times \widetilde{SU}(2) \times U(1)$ symmetry, which remains unchanged in all loop corrections,**

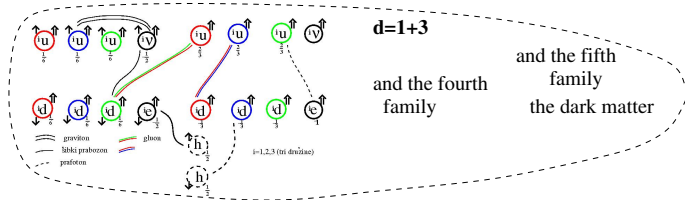
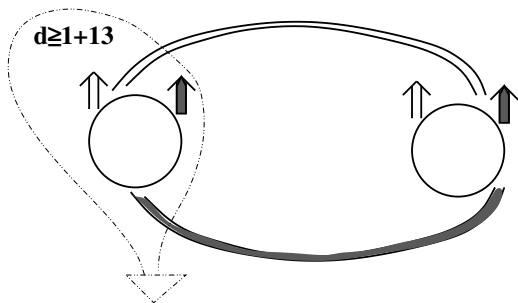
while repetitions of the nonzero vacuum expectation values of the scalar fields, together with loop corrections, both in all orders, give a controllable change of this symmetry.

[arXiv:1902.02691, arXiv:1902.10628]

- ▶ We studied on a toy model of $d = (1 + 5)$ conditions which lead after breaking symmetries to massless spinors chirally coupled to the Kaluza-Klein-like gauge field.,

New J. Phys. **13** (2011) 103027, 1-25,

Int. J Mod. Phys. **A 29**, 1450124 (2014), 21 pages.



**Families of quarks and leptons and
antiquarks and antileptons**

Our technique to represent spinors is elegant.

- ▶ J. of Math. Phys. **34**, 3731 (1993),
- ▶ Int. J. of Modern Phys. **A 9**, 1731 (1994),
- ▶ J. of Math. Phys. **36**, 1593 (1995),
- ▶ J. of Math. Phys. **43**, 5782-5803 (2002), hep-th/0111257,
- ▶ J. of Math. Phys. **44** 4817 (2003), hep-th/0303224,
- ▶ J. of High Ener. Phys. **04** (2014) 165, arxiv:1212.2362v2.

The **spinors** states are created out of
nilpotents $\binom{ab}{\pm i}$ and **projectors** $[\pm i]$

$$\binom{ab}{(\pm i)} := \frac{1}{2}(\gamma^a \mp \gamma^b), \quad [\pm i] := \frac{1}{2}(1 \pm \gamma^a \gamma^b)$$

for $\eta^{aa}\eta^{bb} = -1$,

$$\binom{ab}{(\pm)} := \frac{1}{2}(\gamma^a \pm i\gamma^b), \quad [\pm] := \frac{1}{2}(1 \pm i\gamma^a \gamma^b),$$

for $\eta^{aa}\eta^{bb} = 1$

with γ^a which are usual **Dirac operators** in d -dimensional space.

J. of Math. Phys. **34**, 3731 (1993),

Int. J. of Modern Phys. **A 9**, 1731 (1994),

J. of Math. Phys. **43**, 5782-5803 (2002)

Nilpotents $(\pm i)^{ab}$ and **projectors** $[\pm i]^{ab}$ are eigensates of the **Cartan subalgebra** of S^{ab} and \tilde{S}^{ab} .

$$\begin{aligned} S^{ab}(\mathbf{k}) &= \frac{k^{ab}}{2}(\mathbf{k}), & S^{ab}[\mathbf{k}] &= \frac{k^{ab}}{2}[\mathbf{k}], \\ \tilde{S}^{ab}(\mathbf{k}) &= \frac{k^{ab}}{2}(\mathbf{k}), & \tilde{S}^{ab}[\mathbf{k}] &= -\frac{k^{ab}}{2}[\mathbf{k}]. \end{aligned}$$

γ^a transforms $\binom{ab}{k}$ into $[-k]$, **never** to $\binom{ab}{k}$.

$\tilde{\gamma}^a$ transforms $\binom{ab}{k}$ into $\binom{ab}{k}$, **never** to $[-k]$.

- ▶ One **Weyl representation** of one **family contains** all the **family members** with the **right handed neutrinos included**. It includes also **antimembers**, reachable by either S^{ab} or by $\mathbb{C}_N \mathcal{P}_N$ on a **family member**.

Jour. of High Energy Phys. **04** (2014) 165

- ▶ There are $2^{(7+1)/2-1} = 8$ **families**, which decouple into **twice four families**, with the quantum numbers $(\tilde{\tau}^{2i}, \tilde{N}_R^i)$ and $(\tilde{\tau}^{1i}, \tilde{N}_L^i)$, respectively.

J. of Math. Phys. **34**, 3731 (1993),

Int. J. of Modern Phys. **A 9**, 1731 (1994),

J. of Math. Phys. **44** 4817 (2003), hep-th/030322.,

S^{ab} generate **all the members of one family**. The **eightplet** (represent. of $SO(7, 1)$) of quarks of a particular colour charge

i		$ \psi_i\rangle$	$\Gamma^{(3,1)}$	S^{12}	$\Gamma^{(4)}$	τ^{13}	τ^{23}	Y	τ^4
		Octet, $\Gamma^{(7,1)} = 1$, $\Gamma^{(6)} = -1$, of quarks							
1	u_R^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)(+) & & (+)(+) & & (+) & (-) & (-) \end{matrix}$	1	$\frac{1}{2}$	1	0	$\frac{1}{2}$	$\frac{2}{3}$	$\frac{1}{6}$
2	u_R^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i][-] & & (+)(+) & & (+) & (-) & (-) \end{matrix}$	1	$-\frac{1}{2}$	1	0	$\frac{1}{2}$	$\frac{2}{3}$	$\frac{1}{6}$
3	d_R^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)(+) & & [-][-] & & (+) & (-) & (-) \end{matrix}$	1	$\frac{1}{2}$	1	0	$-\frac{1}{2}$	$-\frac{1}{3}$	$\frac{1}{6}$
4	d_R^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i][-] & & [-][-] & & (+) & (-) & (-) \end{matrix}$	1	$-\frac{1}{2}$	1	0	$-\frac{1}{2}$	$-\frac{1}{3}$	$\frac{1}{6}$
5	d_L^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i](+) & & [-](+) & & (+) & (-) & (-) \end{matrix}$	-1	$\frac{1}{2}$	-1	$-\frac{1}{2}$	0	$\frac{1}{6}$	$\frac{1}{6}$
6	d_L^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)[-] & & [-](+) & & (+) & (-) & (-) \end{matrix}$	-1	$-\frac{1}{2}$	-1	$-\frac{1}{2}$	0	$\frac{1}{6}$	$\frac{1}{6}$
7	u_L^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i](+) & & (+)[-] & & (+) & (-) & (-) \end{matrix}$	-1	$\frac{1}{2}$	-1	$\frac{1}{2}$	0	$\frac{1}{6}$	$\frac{1}{6}$
8	u_L^c	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)[-] & & (+)[-] & & (+) & (-) & (-) \end{matrix}$	-1	$-\frac{1}{2}$	-1	$\frac{1}{2}$	0	$\frac{1}{6}$	$\frac{1}{6}$

$\gamma^0 \gamma^7$ and $\gamma^0 \gamma^8$ transform u_R of the 1st row into u_L of the 7th row, and d_R of the 4th row into d_L of the 6th row, doing what the Higgs scalars and γ^0 do in the Stan. model.

The **anti-octet** (repres. of $SO(7,1)$) of **anti-quarks** of the anti-colour charge, reachable by either S^{ab} or $\mathcal{C}_N \mathcal{P}_N^{(d-1)}$:

i		$ ^a \psi_i \rangle$	$\Gamma^{(3,1)}$	S^{12}	$\Gamma^{(4)}$	τ^{13}	τ^{23}	Y	τ^4
		Antioctet, $\Gamma^{(7,1)} = -1$, $\Gamma^{(6)} = 1$, of antiquarks							
33	$\bar{d}_L^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i](+) & & (+)(+) & & [-] & [+] & [+] \end{matrix}$	-1	$\frac{1}{2}$	1	0	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
34	$\bar{d}_L^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)[-] & & (+)(+) & & [-] & [+] & [+] \end{matrix}$	-1	$-\frac{1}{2}$	1	0	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
35	$\bar{u}_L^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i](+) & & - & & [-] & [+] & [+] \end{matrix}$	-1	$\frac{1}{2}$	1	0	$-\frac{1}{2}$	$-\frac{2}{3}$	$-\frac{1}{6}$
36	$\bar{u}_L^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)[-] & & - & & [-] & [+] & [+] \end{matrix}$	-1	$-\frac{1}{2}$	1	0	$-\frac{1}{2}$	$-\frac{2}{3}$	$-\frac{1}{6}$
37	$\bar{d}_R^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)(+) & & (+)[-] & & [-] & [+] & [+] \end{matrix}$	1	$\frac{1}{2}$	-1	$\frac{1}{2}$	0	$-\frac{1}{6}$	$-\frac{1}{6}$
38	$\bar{d}_R^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i][-] & & (+)[-] & & [-] & [+] & [+] \end{matrix}$	1	$-\frac{1}{2}$	-1	$\frac{1}{2}$	0	$-\frac{1}{6}$	$-\frac{1}{6}$
39	$\bar{u}_R^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ (+i)(+) & & [-](+) & & [-] & [+] & [+] \end{matrix}$	1	$\frac{1}{2}$	-1	$-\frac{1}{2}$	0	$-\frac{1}{6}$	$-\frac{1}{6}$
40	$\bar{u}_R^{\bar{c}1}$	$\begin{matrix} 03 & 12 & 56 & 78 & 9 & 1011 & 1213 & 14 \\ [-i][-] & & [-](+) & & [-] & [+] & [+] \end{matrix}$	1	$-\frac{1}{2}$	-1	$-\frac{1}{2}$	0	$-\frac{1}{6}$	$-\frac{1}{6}$

$\gamma^0 \gamma^7$ and $\gamma^0 \gamma^8$ transform \bar{d}_L of the 1st row into \bar{d}_R of the 5th row, and \bar{u}_L of the 4th row into \bar{u}_R of the 8th row.

Second quantization of fermion fields — in Clifford space manifesting families and family members — of quarks and leptons

- ▶ Defining the creation operator as

$$\hat{b}_1^{1\dagger} = (+i)(+) \mid (+)(+) \parallel (+)(-)(-), \text{ and}$$

- ▶ The annihilation operator as

$$\hat{b}_1^1 = (\hat{b}_1^{1\dagger})^\dagger = (+)(+)(-) \parallel (-)(-) \mid (-)(-i),$$

- ▶ $\hat{b}_2^{1\dagger} = (\hat{b}_2^1)^\dagger = [-i][-] \mid (+)(+) \parallel (+)(-)(-), \text{ and}$

- ▶ $\hat{b}_2^1 = (+)(+)(-) \parallel (-)(-) \mid [-][-i],$

▶

▶

operating on a vacuum state

- ▶ $|\psi_0\rangle = [-i][-][-] \cdots [-] |0\rangle.$

for $d=2(2n+1),$

One finds the commutation relations for a general fermion field which represents just the observed quarks and leptons ($i = (u_{R,L}^\alpha, d_{R,L}^\alpha, \nu_{R,L}^\alpha, e_{R,L}^\alpha)$) in a massless basis and anti-quarks and anti-leptons , with the family quantum numbers α .

- ▶ $\{\hat{\mathbf{b}}_i^\alpha, \hat{\mathbf{b}}_k^{\beta\dagger}\}_+ |\psi_0\rangle = \delta_\beta^\alpha \delta_k^i |\psi_0\rangle ,$
- ▶ $\{\hat{\mathbf{b}}_i^\alpha, \hat{\mathbf{b}}_k^\beta\}_+ |\psi_0\rangle = 0 \cdot |\psi_0\rangle ,$
- ▶ $\{\hat{\mathbf{b}}_i^{\alpha\dagger}, \hat{\mathbf{b}}_k^{\beta\dagger}\}_+ |\psi_0\rangle = 0 \cdot |\psi_0\rangle ,$
- ▶ $\hat{\mathbf{b}}_i^\alpha |\psi_0\rangle = 0 \cdot |\psi_0\rangle ,$
- ▶ $\hat{\mathbf{b}}_i^{\alpha\dagger} |\psi_0\rangle = |\psi_i^\alpha\rangle$

Here the creation and annihilation operators include families. In this case the vacuum state has to be extended.

[arXiv:1802.05554]

In ordinary second quantization procedure the creation operators and correspondingly their Hermitian conjugate operators are **assumed**.

$$\psi_i(\mathbf{t}, \tilde{\mathbf{x}}) = \sum_{\mathbf{p}, i} \hat{b}^\dagger(\mathbf{p}, i) v(\tilde{\mathbf{p}}, i) e^{-\mathbf{p}_a \mathbf{x}^a}.$$

$\hat{b}^\dagger(p, i)$ is just assumed, together with the (assumed) Hermitian conjugate operator), to fill the anticommutation relation.

In the spin-charge-family theory the creation operators appear by themselves.

- ▶ It is worthwhile to notice that "nature could make a choice" of **Grassmann** rather than **Clifford** space:
 - Also in **Grassmann** space, namely, one finds the anticommutation relations needed for a fermion field.
 - But in this case **spinors** would have spins and charges in adjoint representations with respect to particular subgroups.
 - And no **families** would appear.

Vector gauge fields origin in gravity,
in vielbeins and two kinds of the spin connection fields,
the gauge fields of S^{ab} and \tilde{S}^{ab}

- ▶ **All the vector gauge fields, A_m^{Ai} , $(m, n) = (0, 1, 2, 3)$ of the observed charges $\tau^{Ai} = \sum_{s,t} c_{st}^{Ai} S^{st}$, manifesting at the observable energies, **have all the properties as assumed by the standard model.****
- ▶ They carry with respect to the space index $m \in (0, 1, 2, 3)$ the vector degrees of freedom, while they have additional **internal degrees of freedom** (τ^{Ai}) in the adjoint representations.
- ▶ They origin as spin conection gauge fields of S^{ab} :

$$A_m^{Ai} = \sum_{s,t} c^{Aist} \omega_{stm}.$$
- ▶ S^{ab} applies on indexes (s, t, m) as follows

$$S^{ab} \omega_{stm\dots g} = i(\delta_s^a \omega_{tm\dots g}^b - \delta_s^b \omega_{tm\dots g}^a).$$

The action for vectors with respect to the space index
 $m = (0, 1, 2, 3)$ origin in **gravity**

$$\int \mathbf{E} d^4 \mathbf{x} d^{(d-4)} \mathbf{x} \alpha \mathbf{R}^{(d)} = \int d^4 \mathbf{x} \left\{ -\frac{1}{4} \mathbf{F}^{\mathbf{Ai}}{}_{\mathbf{mn}} \mathbf{F}^{\mathbf{Aimn}} \right\},$$

$$\mathbf{A}^{\mathbf{Ai}}{}_{\mathbf{m}} = \sum_{\mathbf{s}, \mathbf{t}} \mathbf{c}^{\mathbf{Aist}} \omega_{\mathbf{stm}}.$$

Eur. Phys. J. C. **77** (2017) 231,

Also scalar fields
(there are doublets and triplets)
origin in gravity fields — **they are spin connections and vielbeins** —
with the space index $s \geq 5$.

Eur. Phys. J. C. **77** (2017) 231 ,
Phys. Rev. **D 91** (2015) 6, 065004,
J. of Mod. Physics **6** (2015) 2244.

- ▶ There are several **scalar gauge fields** with the space index $(s,t,s') = (7,8)$, all origin in the spin connection fields, either $\tilde{\omega}_{abs}$ or $\omega_{s'ts}$:
 - Twice **two triplets**, the scalar gauge fields with the **family** quantum numbers $(\tilde{\tau}^{Ai} = \sum_{a,b} \tilde{c}^{Ai}_{ab} \tilde{S}^{ab})$ and
 - **three singlets** with the **family members** quantum numbers (Q,Q',Y') , the gauge fields of S^{st} .
- ▶ They are all doublets with respect to the space index **(5,6,7,8)**.
- ▶ They have all the rest quantum numbers **determined by the adjoint representations**.
- ▶ They explain at the so far observable energies the **Higgs's scalar** and the **Yukawa couplings**.

The two doublets, determining the properties of the Higgs's scalar and the Yukawa couplings, are:

	state	τ^{13}	$\tau^{23} = Y$	spin	τ^4	Q
A_{78}^{Ai} (-)	$A_7^{Ai} + iA_8^{Ai}$	$+\frac{1}{2}$	$-\frac{1}{2}$	0	0	0
A_{56}^{Ai} (-)	$A_5^{Ai} + iA_6^{Ai}$	$-\frac{1}{2}$	$-\frac{1}{2}$	0	0	-1
A_{78}^{Ai} (+)	$A_7^{Ai} - iA_8^{Ai}$	$-\frac{1}{2}$	$+\frac{1}{2}$	0	0	0
A_{56}^{Ai} (+)	$A_5^{Ai} - iA_6^{Ai}$	$+\frac{1}{2}$	$+\frac{1}{2}$	0	0	+1

There are A_{78}^{Ai} and A_{78}^{Ai} which gain **nonzero vacuum expectation values** at the **electroweak break**.

Index Ai determines the **family** ($\tilde{\tau}^{Ai}$) quantum numbers and the **family members** (Q,Q',Y') quantum numbers, both are in adjoint representations.

- ▶ There are besides **doublets**, with the space index $s = (5, 6, 7, 8)$, as well **triplets** and **anti-triplets**, with respect to the space index $s = (9, \dots, 14)$.
- ▶ **There are no additional scalars** in the theory for **$d=(13+1)$** .
- ▶ **All are massless.**
- ▶ All the scalars have the family and the family members quantum numbers in the **adjoint** representations.
- ▶ The properties of scalars are to be analyzed with respect to the generators of the corresponding subgroups, expressible with S^{ab} , as it is the case of the vector gauge fields.
- ▶ It is the (**so far assumed**) **condensate**, which makes those gauge fields, with which it interacts, massive.
 - o **The condensate breaks the CP symmetry.**

- ▶ The **scalar condensate** of two **right handed neutrinos** couple to
 - all the **scalar and vector** gauge fields, making them massive,
 - It does not interact with the **weak charge $SU(2)_I$** , the **hyper charge $U(1)$** , and the **colour $SU(3)$ charge gauge fields**, as well as the **gravity**, leaving them **massless**.

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Phys Rev.**D 91**(2015)6,065004.

Scalars with $\mathbf{s}=(7,8)$, which gain **nonzero vacuum expectation values**, break the **weak and the hyper** symmetry, while conserving the **electromagnetic and colour** charge:

$$\begin{aligned} \mathbf{A}_s^{\mathbf{A}i} &\supset (\mathbf{A}_s^{\mathbf{Q}}, \mathbf{A}_s^{\mathbf{Q}'}, \mathbf{A}_s^{\mathbf{Y}'}, \tilde{\tilde{\mathbf{A}}}_s^{\tilde{\tilde{1}}}, \tilde{\tilde{\mathbf{A}}}_s^{\tilde{\tilde{N}}_L}, \tilde{\tilde{\mathbf{A}}}_s^{\tilde{\tilde{2}}}, \tilde{\tilde{\mathbf{A}}}_s^{\tilde{\tilde{N}}_R}), \\ \tau^{\mathbf{A}i} &\supset (\mathbf{Q}, \mathbf{Q}', \mathbf{Y}', \tilde{\tilde{\tau}}^1, \tilde{\tilde{\mathbf{N}}}_L, \tilde{\tilde{\tau}}^2, \tilde{\tilde{\mathbf{N}}}_R), \\ \mathbf{s} &= (7, 8). \end{aligned}$$

Ai denotes:

o family quantum numbers

$(\tilde{\tilde{\tau}}^1, \tilde{\tilde{\mathbf{N}}}_L)$ quantum numbers of the first group of four families

and

$(\tilde{\tilde{\tau}}^2, \tilde{\tilde{\mathbf{N}}}_R)$ quantum numbers of the second group of four families.

o And family members quantum numbers $(\mathbf{Q}, \mathbf{Q}', \mathbf{Y}')$

A_s^{Ai} are expressible with either $\omega_{sts'}$ or $\tilde{\omega}_{abs'}$.

$$\vec{A}_s^1 = (\tilde{\omega}_{58s} - \tilde{\omega}_{67s}, \tilde{\omega}_{57s} + \tilde{\omega}_{68s}, \tilde{\omega}_{56s} - \tilde{\omega}_{78s}),$$

$$\vec{A}_s^2 = (\tilde{\omega}_{58s} + \tilde{\omega}_{67s}, \tilde{\omega}_{57s} - \tilde{\omega}_{68s}, \tilde{\omega}_{56s} + \tilde{\omega}_{78s}),$$

$$\vec{A}_{Ls}^N = (\tilde{\omega}_{23s} + i\tilde{\omega}_{01s}, \tilde{\omega}_{31s} + i\tilde{\omega}_{02s}, \tilde{\omega}_{12s} + \tilde{\omega}_{03s}),$$

$$\vec{A}_{Rs}^N = (\tilde{\omega}_{23s} - i\tilde{\omega}_{01s}, \tilde{\omega}_{31s} - i\tilde{\omega}_{02s}, \tilde{\omega}_{12s} - i\tilde{\omega}_{03s}),$$

$$A_s^Q = \omega_{56s} - (\omega_{910s} + \omega_{1112s} + \omega_{1314s}),$$

$$A_s^Y = (\omega_{56s} + \omega_{78s}) - (\omega_{910s} + \omega_{1112s} + \omega_{1314s})$$

$$A_s^4 = -(\omega_{910s} + \omega_{1112s} + \omega_{1314s}).$$

The **mass term**, appearing in the **starting action**,
 is (p_s , when treating the lowest energy solutions, is left out)

$$\mathcal{L}_M = \sum_{s=(7,8), Ai} \bar{\psi} \gamma^s (-\tau^{Ai} A_s^{Ai}) \psi =$$

$$-\bar{\psi} \left\{ \overset{78}{(+)} \tau^{Ai} (A_7^{Ai} - i A_8^{Ai}) + \overset{78}{(-)} \tau^{Ai} (A_7^{Ai} + i A_8^{Ai}) \right\} \psi ,$$

$$\overset{78}{(\pm)} = \frac{1}{2} (\gamma^7 \pm i \gamma^8), \quad A_{78(\pm)}^{Ai} := (A_7^{Ai} \mp i A_8^{Ai}).$$

Operators Y , Q and τ^{13} , applied on $(A_7^{A_i} \mp i A_8^{A_i})$

$$\tau^{13} (A_7^{A_i} \mp i A_8^{A_i}) = \pm \frac{1}{2} (A_7^{A_i} \mp i A_8^{A_i}),$$

$$Y (A_7^{A_i} \mp i A_8^{A_i}) = \mp \frac{1}{2} (A_7^{A_i} \mp i A_8^{A_i}),$$

$$Q (A_7^{A_i} \mp i A_8^{A_i}) = 0,$$

manifest that **all** $(A_7^{A_i} \mp i A_8^{A_i})$ have quantum numbers of the **Higgs's scalar of the standard model**, "dressing", after **gaining nonzero expectation values**, the right handed members of a family with appropriate charges, so that they gain charges of the left handed partners:

$(A_7^{A_i} + i A_8^{A_i})$ "dresses" u_R, ν_R and $(A_7^{A_i} - i A_8^{A_i})$ "dresses" d_R, e_R , with quantum numbers of their left handed partners, just as required by the "standard model".

Ai measures:

either

o the **Q,Q',Y'** charges of the **family members**

or

o **family** charges ($\vec{\tau}^{\vec{I}}, \vec{N}_L$), transforming a family member of one family into the same family member of another family,

manifesting in each group of four families the

$$\widetilde{SU}(2) \times \widetilde{SU}(2) \times U(1)$$

symmetry.

Eight families of u_R (spin 1/2, colour $(\frac{1}{2}, \frac{1}{2\sqrt{3}})$) and of colourless ν_R (spin 1/2). All have "tilde spinor charge" $\tilde{\tau}^4 = -\frac{1}{2}$, the weak charge $\tau^{13} = 0$, $\tau^{23} = \frac{1}{2}$. Quarks have "spinor" q.no. $\tau^4 = \frac{1}{6}$ and leptons $\tau^4 = -\frac{1}{2}$. The first four families have $\tilde{\tau}^{23} = 0$, $\tilde{N}_R^3 = 0$, the second four families have $\tilde{\tau}^{13} = 0$, $\tilde{N}_L^3 = 0$.

$\tilde{N}_R^3 = 0, \tilde{\tau}^{23} = 0$				$\tilde{N}_R^3 = 0, \tilde{\tau}^{23} = 0$				$\tilde{\tau}^{13}$	\tilde{N}_L^3
u_{R1}^{c1}	03 12 (+i) [+] [+] (+) (+) [-] [-]	56 78 [+] (+) [+] (+) (+) [-] [-]	9 10 11 12 13 14	ν_{R1}	03 12 (+i) [+] [+] (+) (+) (+) (+)	56 78 [+] (+) [+] (+) (+) (+) (+)	9 10 11 12 13 14	$-\frac{1}{2}$	$-\frac{1}{2}$
u_{R2}^{c1}	03 12 [+ i] (+) [+] (+) (+) [-] [-]	56 78 [+] (+) [+] (+) (+) [-] [-]	9 10 11 12 13 14	ν_{R2}	03 12 [+ i] (+) [+] (+) (+) (+) (+)	56 78 [+] (+) [+] (+) (+) (+) (+)	9 10 11 12 13 14	$-\frac{1}{2}$	$\frac{1}{2}$
u_{R3}^{c1}	03 12 (+i) [+] (+) [+] (+) [-] [-]	56 78 (+) [+] (+) [+] (+) [-] [-]	9 10 11 12 13 14	ν_{R3}	03 12 (+i) [+] (+) [+] (+) (+) (+)	56 78 (+) [+] (+) [+] (+) (+) (+)	9 10 11 12 13 14	$\frac{1}{2}$	$-\frac{1}{2}$
u_{R4}^{c1}	03 12 [+ i] (+) (+) [+] (+) [-] [-]	56 78 [+] (+) (+) [+] (+) [-] [-]	9 10 11 12 13 14	ν_{R4}	03 12 [+ i] (+) (+) [+] (+) (+) (+)	56 78 [+] (+) (+) [+] (+) (+) (+)	9 10 11 12 13 14	$\frac{1}{2}$	$\frac{1}{2}$
$\tilde{N}_L^3 = 0, \tilde{\tau}^{13} = 0$				$\tilde{N}_L^3 = 0, \tilde{\tau}^{13} = 0$				$\tilde{\tau}^{23}$	\tilde{N}_R^3
u_{R5}^{c1}	03 12 (+i) (+) (+) (+) (+) [-] [-]	56 78 (+) (+) (+) (+) (+) [-] [-]	9 10 11 12 13 14	ν_{R5}	03 12 (+i) (+) (+) (+) (+) (+) (+)	56 78 (+) (+) (+) (+) (+) (+) (+)	9 10 11 12 13 14	$-\frac{1}{2}$	$-\frac{1}{2}$
u_{R6}^{c1}	03 12 (+i) (+) [+] [+] (+) [-] [-]	56 78 (+) (+) [+] [+] (+) [-] [-]	9 10 11 12 13 14	ν_{R6}	03 12 (+i) (+) [+] [+] (+) (+) (+)	56 78 (+) (+) [+] [+] (+) (+) (+)	9 10 11 12 13 14	$-\frac{1}{2}$	$\frac{1}{2}$
u_{R7}^{c1}	03 12 [+ i] [+] (+) (+) (+) [-] [-]	56 78 (+) (+) (+) (+) (+) [-] [-]	9 10 11 12 13 14	ν_{R7}	03 12 [+ i] [+] (+) (+) (+) (+) (+)	56 78 (+) (+) (+) (+) (+) (+) (+)	9 10 11 12 13 14	$\frac{1}{2}$	$-\frac{1}{2}$
u_{R8}^{c1}	03 12 [+ i] [+] [+] [+] (+) [-] [-]	56 78 [+] [+] [+] [+] (+) [-] [-]	9 10 11 12 13 14	ν_{R8}	03 12 [+ i] [+] [+] [+] (+) (+) (+)	56 78 [+] [+] [+] [+] (+) (+) (+)	9 10 11 12 13 14	$\frac{1}{2}$	$\frac{1}{2}$

Before the **electroweak break** all the **families** are **mass protected** and correspondingly **massless**.

- ▶ Scalars with the weak and the hyper charge $(\mp\frac{1}{2}, \pm\frac{1}{2})$ determine masses of **all the family members α** of the **lower four families**, ν_R of the lower four families have nonzero $Y' := -\tau^4 + \tau^{23}$ and interact with the scalar field $(A_{(\pm)}^{Y'}, \vec{A}_{(\pm)}^{\tilde{1}}, \vec{A}_{(\pm)}^{\tilde{N}_L})$.
- ▶ The group of the lower four families manifest the $\widetilde{SU}(2)_{\widetilde{SO}(1,3)} \times \widetilde{SU}(2)_{\widetilde{SO}(4)} \times U(1)$ **symmetry** (also after all loop corrections).

$$\mathcal{M}^\alpha = \begin{pmatrix} -a_1 - a & e & d & b \\ e^* & -a_2 - a & b & d \\ d^* & b^* & a_2 - a & e \\ b^* & d^* & e^* & a_1 - a \end{pmatrix}^\alpha.$$

[arXiv:1412.5866], [arXiv:1902.02691], [arXiv:1902.10628]

We **made calculations**, treating **quarks** and **leptons** in equivalent way, as required by the "spin-charge-family" theory. Although

- ▶ any **$(n-1) \times (n-1)$** submatrix of an unitary **$n \times n$** matrix determines the **$n \times n$** matrix for **$n \geq 4$** uniquely,
- ▶ the **measured mixing matrix elements** of the **3×3** submatrix **are not yet accurate enough even for quarks to predict the masses m_4 of the fourth family members.**
 - We can say, taking into account the data for the mixing matrices and masses, that **m_4 quark masses might be any in the interval $(300 < m_4 < 1000)$ GeV or even **above**.** Other experiments require that m_4 are above 1000 GeV.
- ▶ **Assuming** masses **m_4** we can predict mixing matrices.

Results are presented for two choices of $m_{u_4} = m_{d_4}$,
 [arxiv:1412.5866]:

- ▶ 1. $m_{u_4} = 700$ GeV, $m_{d_4} = 700$ GeV.....new₁
- ▶ 2. $m_{u_4} = 1\,200$ GeV, $m_{d_4} = 1\,200$ GeV.....new₂

$$|V_{(ud)}| = \begin{pmatrix} \begin{array}{cccc} \text{exp}_n & 0.97425 \pm 0.00022 & 0.2253 \pm 0.0008 & 0.00413 \pm 0.00049 \\ \text{new}_1 & 0.97423(4) & 0.22539(7) & 0.00299 & \mathbf{0.00776(1)} \\ \text{new}_2 & 0.97423[5] & 0.22538[42] & 0.00299 & \mathbf{0.00793[466]} \end{array} \\ \begin{array}{cccc} \text{exp}_n & 0.225 \pm 0.008 & 0.986 \pm 0.016 & 0.0411 \pm 0.0013 \\ \text{new}_1 & 0.22534(3) & 0.97335 & 0.04245(6) & \mathbf{0.00349(60)} \\ \text{new}_2 & 0.22531[5] & 0.97336[5] & 0.04248 & \mathbf{0.00002[216]} \end{array} \\ \begin{array}{cccc} \text{exp}_n & 0.0084 \pm 0.0006 & 0.0400 \pm 0.0027 & 1.021 \pm 0.032 \\ \text{new}_1 & 0.00667(6) & 0.04203(4) & 0.99909 & \mathbf{0.00038} \\ \text{new}_2 & 0.00667 & 0.04206[5] & 0.99909 & \mathbf{0.00024[21]} \end{array} \\ \begin{array}{cccc} \text{new}_1 & \mathbf{0.00677(60)} & \mathbf{0.00517(26)} & \mathbf{0.00020} & \mathbf{0.99996} \\ \text{new}_2 & \mathbf{0.00773} & \mathbf{0.00178} & \mathbf{0.00022} & \mathbf{0.99997[9]} \end{array} \end{pmatrix} .$$

One can see what

B. Belfatto, R. Beradze, Z. Berezhiani, required in
 [arXiv:1906.02714v1], that

$V_{u_1 d_4} > V_{u_1 d_3}$, $V_{u_2 d_4} < V_{u_1 d_4}$, and $V_{u_3 d_4} < V_{u_1 d_4}$,
 what is just happening in my theory.

- ▶ ○ The **matrix elements** V_{CKM} **depend strongly on the accuracy** of the experimental **3 x 3 submatrix**.
- Calculated **3 x 3 submatrix** of 4 x 4 V_{CKM} depends on the m_{4th} **family masses**, but not much.
- $V_{u_i d_4}$, $V_{d_j u_4}$ do not depend strongly on the m_{4th} family masses and are obviously **very small**.
- ▶ The higher are the fourth family members masses, the closer are the mass matrices to the **democratic matrices** for either quarks or leptons, as expected.
- ▶ The higher are the fourth family members masses, the better are conditions

$$V_{u_1 d_4} > V_{u_1 d_3} ,$$

$$V_{u_2 d_4} < V_{u_1 d_4} , \text{ and}$$

$$V_{u_3 d_4} < V_{u_1 d_4}$$

fulfilled.

- ▶ The **stable family** of the **upper four families** group is the candidate to form the **Dark Matter**.
- ▶ Masses of the upper four families are influenced :
 - by the $\widetilde{SU}(2)_{\parallel \widetilde{SO}(3,1)} \times \widetilde{SU}(2)_{\parallel \widetilde{SO}(4)}$ **scalar fields** with the corresponding family quantum numbers,
 - by the **scalars** $(A_{78}^Q, A_{78}^{Q'}, A_{78}^{Y'})$, and
 - by the **condensate** of the two ν_R of the **upper four families**.

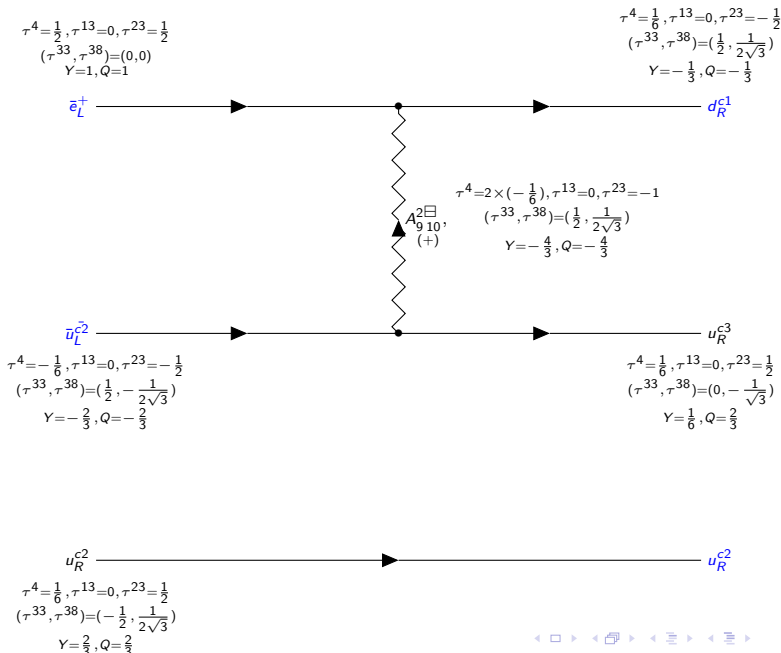
Matter-antimatter asymmetry

There are also **triplet** and **anti-triplet** scalars, $s = (9, \dots, d)$:

	state	τ^{33}	τ^{38}	spin	τ^4	Q
$A_{9\ 10}^{Ai}$ (+)	$A_9^{Ai} - iA_{10}^{Ai}$	$+\frac{1}{2}$	$\frac{1}{2\sqrt{3}}$	0	$-\frac{1}{3}$	$-\frac{1}{3}$
$A_{11\ 12}^{Ai}$ (+)	$A_{11}^{Ai} - iA_{12}^{Ai}$	$-\frac{1}{2}$	$\frac{1}{2\sqrt{3}}$	0	$-\frac{1}{3}$	$-\frac{1}{3}$
$A_{13\ 14}^{Ai}$ (+)	$A_{13}^{Ai} - iA_{14}^{Ai}$	0	$-\frac{1}{\sqrt{3}}$	0	$-\frac{1}{3}$	$-\frac{1}{3}$
$A_{9\ 10}^{Ai}$ (-)	$A_9^{Ai} + iA_{10}^{Ai}$	$-\frac{1}{2}$	$-\frac{1}{2\sqrt{3}}$	0	$+\frac{1}{3}$	$+\frac{1}{3}$
$A_{11\ 12}^{Ai}$ (-)	$A_{11}^{Ai} + iA_{12}^{Ai}$	$\frac{1}{2}$	$-\frac{1}{2\sqrt{3}}$	0	$+\frac{1}{3}$	$+\frac{1}{3}$
$A_{13\ 14}^{Ai}$ (-)	$A_{13}^{Ai} + iA_{14}^{Ai}$	0	$\frac{1}{\sqrt{3}}$	0	$+\frac{1}{3}$	$+\frac{1}{3}$

They cause transitions from anti-leptons into quarks and anti-quarks into quarks and back, **transforming matter into antimatter and back**. The condensate breaks CP symmetry, offering the explanation for the matter-antimatter asymmetry in the universe.

Let us look at scalar triplets, causing the birth of a proton from the left handed **positron**, **antiquark** and **quark**:



These two quarks, d_R^{c1} and u_R^{c3} can bind (at low enough energy) together with u_R^{c2} into the colour **chargeless baryon - a proton**.

After the appearance of the **condensate** the **CP is broken**.

In the expanding universe, fulfilling the Sakharov request for appropriate non-thermal equilibrium, **these triplet scalars have a chance to explain the matter-antimatter asymmetry**.

The opposite transition makes the proton decay.

These processes seem to explain the lepton number non conservation.

Dark matter

$d \rightarrow (d - 4) + (3 + 1)$ before (or at least at) the electroweak break.

- ▶ We follow the **evolution of the universe**, in particular the **abundance of the fifth family members** - the **candidates** for the **dark matter** in the universe.
- ▶ We estimate the behaviour of our stable heavy family quarks and anti-quarks in the expanding universe by solving the system of **Boltzmann equations**.
- ▶ We follow the **clustering** of the **fifth family** quarks and antiquarks into the **fifth family baryons** through the **colour** phase transition.
- ▶ The **mass** of the fifth family members is determined from the today **dark matter density**.

Phys. Rev. D (2009) 80.083534

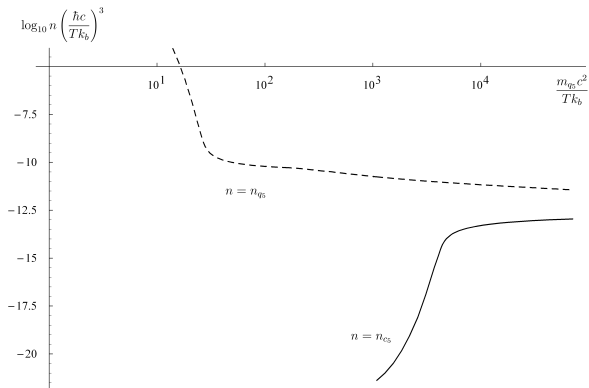


Figure: The dependence of the two number densities n_{q_5} (of the fifth family quarks) and n_{c_5} (of the fifth family clusters) as the function of $\frac{m_{q_5} c^2}{T k_b}$ is presented for the values $m_{q_5} c^2 = 71 \text{ TeV}$, $\eta_{c_5} = \frac{1}{50}$ and $\eta_{(q\bar{q})_b} = 1$. We take $g^* = 91.5$.

We estimated from following the fifth family members in the expanding universe:



$$10 \text{ TeV} < m_{q_5} c^2 < 4 \cdot 10^2 \text{ TeV} .$$



$$10^{-8} \text{ fm}^2 < \sigma_{c_5} < 10^{-6} \text{ fm}^2 .$$

(It is at least $10^{-6} \times$ smaller than the cross section for the first family neutrons.)

We estimate from the scattering of the fifth family members on the ordinary matter on our Earth, on the direct measurements - DAMA, CDMS,...- ...



$$200 \text{ TeV} < m_{q_5} c^2 < 10^5 \text{ TeV} .$$

Second quantization in Clifford and in Grassmann space.

In ordinary second quantization procedure the creation operators and correspondingly their Hermitian conjugate operators are **assumed**.

$$\psi_i(\mathbf{t}, \tilde{\mathbf{x}}) = \sum_{\mathbf{p}, i} \hat{b}^\dagger(\mathbf{p}, i) v(\tilde{\mathbf{p}}, i) e^{-\mathbf{p}_a \mathbf{x}^a}.$$

$\hat{b}^\dagger(p, i)$ is just assumed (together with the (assumed) Hermitian conjugate operator) to fulfill the anticommutation relation.

In the spin-charge-family theory the creation operators appear **by themselves**.

Let us recognize that the **description of the internal degrees of freedom in Clifford space** do not need any assumption about the creation and annihilation operators

- ▶ In $d = 2(2n + 1)$ from the starting state made as a product of an odd number of only nilpotents, for example, $|\psi_1^1\rangle = \hat{b}_1^{1\dagger} |\psi_{oc}\rangle$,

$$\hat{b}_1^{1\dagger} : = (+i)(+)(+) \cdots (+) (+),$$

$$\hat{b}_1^1 = (\hat{b}_1^{1\dagger})^\dagger = (-) (-) \cdots (-)(-)(-i),$$

- ▶ all the other creation operators and correspondingly states, follow by the application of products of S^{ab} and $\tilde{\zeta}^{ab}$

$$\begin{aligned} \hat{b}_i^{\alpha\dagger} &\propto \tilde{\zeta}^{ab} \dots \tilde{\zeta}^{ef} S^{mn} \dots S^{pr} \hat{b}_1^{1\dagger} \\ &\propto S^{mn} \dots S^{pr} \hat{b}_1^{1\dagger} S^{ab} \dots S^{ef}. \end{aligned}$$

$$\hat{b}_i^\alpha = (\hat{b}_i^{\alpha\dagger})^\dagger \propto S^{ef} \dots S^{ab} \hat{b}_1^1 S^{pr} \dots S^{mn}.$$

These $2^{\frac{d}{2}-1} \times 2^{\frac{d}{2}-1}$ creation operators and the same number of annihilation operators obey the anti-commutation relations required for fermions.

$$\begin{aligned} \{\hat{b}_i^\alpha, \hat{b}_j^{\alpha'\dagger}\}_+ |\psi_{oc}\rangle &= \delta^{\alpha\alpha'} \delta_{ij} |\psi_{oc}\rangle, \\ \{\hat{b}_i^\alpha, \hat{b}_j^{\alpha'}\}_+ |\psi_{oc}\rangle &= 0 |\psi_{oc}\rangle, \\ \{\hat{b}_i^{\alpha\dagger}, \hat{b}_j^{\alpha'\dagger}\}_+ |\psi_{oc}\rangle &= 0 |\psi_{oc}\rangle, \\ \hat{b}_i^\alpha |\psi_{oc}\rangle &= 0 |\psi_{oc}\rangle, \\ \hat{b}_i^{\alpha\dagger} |\psi_{oc}\rangle &= |\psi_i^\alpha\rangle, \end{aligned}$$

with (i, j) determining family members quantum numbers and (α, α') denoting "family" quantum numbers.

$$|\psi_{oc}\rangle = \begin{matrix} 03 & 12 & 56 & & d-1 & d & & 03 & 12 & 56 & & d-1 & d & & 03 & 12 & 56 \\ [-i] & [-] & [-] & \cdots & [-] & & & [+i] & [+] & [-] & \cdots & [-] & & & [+i] & [-] & [+] & \cdots \end{matrix}$$

for $d=2(2n+1)$.

For a n particle state of fermions and anti-fermions above the "Dirac sea" $|1_{sp_1}^\alpha, 1_{s'p_2}^{\alpha'}, 1_{s''p_3}^{\alpha''}, \dots, 0_{s'''p_i}^{\alpha'''}, \dots, 0_{s^{iv}p_j}^{\alpha^{iv}}, \dots, \rangle$
it follows



$$\begin{aligned} & \hat{b}_{s'''p_i}^{\alpha''' \dagger} \hat{b}_{s^{iv}p_j}^{\alpha^{iv} \dagger} |1_{sp_1}^\alpha, 1_{s'p_2}^{\alpha'}, 1_{s''p_3}^{\alpha''}, \dots, 0_{s'''p_i}^{\alpha'''}, \dots, 0_{s^{iv}p_j}^{\alpha^{iv}}, \dots, \rangle = \\ - & \hat{b}_{s^{iv}p_j}^{\alpha^{iv} \dagger} \hat{b}_{s'''p_i}^{\alpha''' \dagger} |1_{sp_1}^\alpha, 1_{s'p_2}^{\alpha'}, 1_{s''p_3}^{\alpha''}, \dots, 0_{s'''p_i}^{\alpha'''}, \dots, 0_{s^{iv}p_j}^{\alpha^{iv}}, \dots, \rangle, \end{aligned}$$

where 1 denotes the occupied state and 0 the unoccupied state.



$$\begin{aligned} \hat{N}_{sp_k}^\alpha &= \hat{b}_{sp_k}^{\alpha \dagger} \hat{b}_{sp_k}^\alpha, \quad (\hat{N}_{sp_k}^\alpha)^2 = \hat{N}_{sp_k}^\alpha, \\ \hat{N} &= \sum_{\alpha, s, p_k} \hat{N}_{sp_k}^\alpha. \end{aligned}$$

The number operator $\hat{N}_{sp_k}^\alpha$ has the eigenvalues 1 or 0.

- ▶ **Each single particle state carries its own internal space**, described by a creation operator with a superposition of an odd number of γ_i^a 's, and **its own coordinate space**, described by x_i^a 's (or p_i^a).
The **creation operators of any two pairs of particles therefore anti-commute**.
- ▶ Correspondingly the **two states of two fermions must distinguish in either internal space or in the coordinate space**.
- ▶ The property of the creation operators $\hat{b}_{spi}^{\alpha\dagger} \hat{b}_{s'p'j}^{\alpha'\dagger}$ applying on the n-particle state $|1_{sp1}^{\alpha}, 1_{s'p'2}^{\alpha'}, 1_{s''p''3}^{\alpha''}, \dots, 0_{s'''p_i}^{\alpha'''}, \dots, 0_{s^{iv}p_j}^{\alpha^{iv}}, \dots, \rangle$, can be as well described by (superposition of) Slater determinants of single particle states.

- ▶ In the **Clifford** case all the spins and the charges and the family quantum numbers belong to the fundamental representations of the groups.
- ▶ I demonstrated to you that **Clifford space does not need to postulate creation and annihilation operators.**
- ▶ Let me add that also the **Grassmann space does not need to postulate creation and annihilation operators.**
- ▶ **The creation and annihilation operators in Grassmann space as well anticommute as it is required for the fermions!**
- ▶ And yet their **spins are integer and their charges correspondingly in the adjoint representations.**
- ▶ **Could there exist fermions with integer spins instead of with the half integer spins?**

- ▶ In the *standard model* the **family members** with all their properties, the **families**, the **gauge vector fields**, the **scalar Higgs**, the **Yukawa couplings**, exist by the **assumption**.
- ▶ ** In the **spin-charge-family theory** the appearance and all the properties of all these fields follow from the simple starting action with **two kinds of spins** and with the **gravity only** .
 - ** The theory offers the explanation for the **dark matter**.
 - ** The theory offers the explanation for the **matter-antimatter asymmetry**.
 - ** All the **scalar** and all the **vector** gauge fields are **directly or indirectly observable**.
- ▶ ** The **spin-charge-family theory** even offers the **creation and annihilation operators without postulation**.

The *spin-charge-family theory* explains also many other properties, which are not explainable in the *standard model*, like "miraculous" non-anomalous triangle Feynman diagrams.

The more work is put into the *spin-charge-family theory* the more explanations for the phenomena follow.

Concrete predictions:

- ▶ There are several scalar fields;
 - **two triplets** , ◦ **three singlets** ,explaining **higgs** and **Yukawa couplings**, some of them will be observed at the LHC, JMP 6 (2015) 2244,
Phys. Rev. D 91 (2015) 6, 065004.
- ▶ There is the **fourth family**, (weakly) coupled to the observed **three**, which will be observed at the LHC, New J. of Phys. 10 (2008) 093002.
- ▶ There is the **dark matter** with the predicted properties, Phys. Rev. D (2009) 80.083534.
- ▶ There is the ordinary **matter/antimatter asymmetry** explained and the **proton decay** predicted and explained, Phys. Rev. D 91 (2015) 6, 065004.

We recognize that:

- ▶ The last **data for mixing matrix of quarks** are in better agreement with our prediction for the **3×3 submatrix** elements of the **4×4 mixing matrix** than the previous ones.
- ▶ Our **fit** to the last data predicts how will the **3×3 submatrix elements change** in the next more accurate measurements.
- ▶ Masses of the **fourth family lie much above** the known three, masses of quarks are close to each other.
- ▶ The larger are **masses of the fourth family the larger are $V_{u_1 d_4}$ in comparison with $V_{u_1 d_3}$ and the more is valid that $V_{u_2 d_4} < V_{u_1 d_4}$, $V_{u_3 d_4} < V_{u_1 d_4}$.**
The flavour changing neutral currents are correspondingly weaker.

- ▶ Masses of the **fifth family** lie **much above** the known three and the **predicted fourth family** masses.
- ▶ Although the upper four families carry the weak (of two kinds) and the colour charge, these group of four families are completely decoupled from the lower four families up to the $< 10^{16}$ GeV, unless the breaks of symmetries recover.
- ▶ **Baryons** of the **fifth family** are heavy, forming small enough clusters with small enough scattering amplitude among themselves and with the ordinary matter to be the candidate for the **dark matter**.
- ▶ The "nuclear" force among them is different from the force among ordinary nucleons.

- ▶ The **spin-charge-family** theory is offering an explanation for the **hierarchy problem**:
The mass matrices of the **two four families groups** are almost democratic, causing spreading of the **fermion masses** from 10^{16} GeV to 10^{-8} MeV.

To summarize:

- ▶ I hope that I managed to convince you that I can answer many open questions of particle physics and cosmology.
- ▶ **The collaborators are very welcome!**
- ▶ There are namely a lot of properties to derive.